Postannealing effects on magnetic properties and microstructure of CoCrPt/Ti perpendicular recording media

Anup G. Roy
Data Storage Systems Center, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213-3890

N. T. Nuhfer
Department of Materials Science and Engineering, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213-3890

David E. Laughlin
Data Storage Systems Center and Department of Materials Science and Engineering, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213-3890

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In this work, we investigate the postannealing effects on the magnetic and structural properties of CoCrPt perpendicular films. We observe a coercivity of 5000 Oe in the films with a 2 nm CrMn overlayer, which is about two times larger than the coercivity of similar films without a CrMn overlayer. This increment is attributed to the decoupling of grains by diffusion of CrMn from the top layer through the grain boundaries. An increase in the negative nucleation field and a decrease in intergranular exchange coupling with annealing temperatures was observed for the films with a CrMn overlayer. On the other hand, the films without a CrMn overlayer show the opposite trends except at high annealing temperature (450 °C). We observe a coercivity of ~7600 Oe and a negative nucleation field of ~2400 Oe for a film with a CrMn overlayer annealed at 450 °C for 5 min.

High coercivity, squareness, and negative nucleation field are very important for perpendicular magnetic recording media. The current challenge is to improve these properties. It is generally agreed that segregation of nonmagnetic elements to the boundaries of the columnar grains can lead to higher coercivity and lower media noise. Also, substrate heating and postannealing are the common practices used to improve magnetic properties for Co based longitudinal media by segregation and interdiffusion to the grain boundaries of nonmagnetic elements. By such interdiffusion both Cr or/ and Mn have been shown to increase the coercivity of Co based longitudinal media.1–4 The effect of postannealing on Co based perpendicular media has yet to be investigated. This article presents the effect of Cr and Mn interdiffusion by rapid thermal postannealing (RTA) on the magnetic properties and microstructure of CoCrPt perpendicular media.

Two types of CoCrPt thin films were deposited at about 280 °C by rf diode sputtering onto heated naturally oxidized Si (111) wafers: Type I (without a CrMn layer): Si substrate/Ti (50 nm)/CoCrPt (30 nm) and type II (with a CrMn layer): Si substrate/Ti (50 nm)/CoCrPt (30 nm)/CrMn (2 nm). Postdeposition rapid thermal annealing was performed on the samples at atmospheric pressure under Ar flow. The Ar sputtering gas pressure was 5 mTorr and the base pressure was below 5 × 10⁻⁷ Torr. The sputtering rate was about 0.1 nm/s. The magnetic properties were measured using an alternating field detectable in both hysteresis loops is due to an initial non-textured layer in the CoCrPt films as detected by high-resolution transmission electron microscopy (TEM).5

Typical loops at different annealing temperatures for films without CrMn layer (type I) are shown in Fig. 2(a). Once again the loops are different. The slope (dM/dH) at Hc became steeper with anneal temperature. However, the slope decreases for the loop of 450 °C. The coercivity and negative nucleation field decrease with annealing temperatures up to 400 °C and then they increase slightly at 450 °C, as shown in...
Fig. 3 II films at different annealing temperatures are shown in the two different types of films. The plots for type I and type II films show the opposite trend than that of type I films. The loops became more and more slanted with increasing annealing temperatures. The coercivity and the negative nucleation field increase monotonously with annealing temperature for type II films [Fig. 2(c)]. (\(H_c\) and \(-H_n\) values should be higher than the measured value since the saturation field cannot be reached for the films annealed at 400 and 450 °C due to the field limitation of the apparatus.) The improvement of the magnetic properties is due to the interdiffusion of Cr and Mn into the grain boundaries of the CoCrPt layer from the top layer as discussed later.

The exchange coupling in the film can be qualitatively understood from the differences of the value of switching volume (\(V^*\)). \(V^*\) can be evaluated by the measurement of the dependence of the \(H_c\) on the sweep rate of the applied field. Since the presence of the demagnetizing field causes the internal field to change following the variation of the magnetization during the measurement, the measurement of the variation of the magnetization with time in the presence of a steady negative field (time decay method) to determine the \(H_c\) cannot be applied to the films with perpendicular anisotropy. Here we have measured the coercivities with various sweep rates (13–1300 Oe/s) to determine the \(V^*\) for the two different types of films. The plots for type I and type II films at different annealing temperatures are shown in Figs. 3(a) and 3(b), respectively. As seen from the figures, the \(H_c\) displays a linear relationship with the logarithm of the sweep rate. From the slope of the plot, the switching volumes can be calculated. The calculated \(V^*\) for both types of film are plotted as a function of annealing temperatures [Fig. 3(c)]. \(V^*\) is an estimate of volume that is coherently changing the moment due to the applied field. The plot for type I films in Fig. 3(c) shows that the \(V^*\) value rises with annealing temperatures up to 400 °C. This indicates an increase of exchange coupling in the films with annealing time. The relatively lower \(V^*\) value for the film annealed at 450 °C indicates that the film becomes relatively less exchange coupled by the segregation of Cr to the grain boundaries at this higher temperature. On the other hand, the plot for type II films shows that the \(V^*\) value falls linearly with annealing temperatures. This trend of the plot implies that the type II films become more exchanged decoupled due to diffusion of more Cr and Mn into the grain boundaries from the top layer with anneal-temperature.

To observe any microstructural changes in our film due to annealing, we have performed plan-view and cross-sectional TEM. Figures 4(a) and 4(b) show typical plan-view bright-field images of as-deposited and annealed at 450 °C type I films, respectively. Figures 4(c) and 4(d) are as-deposited and annealed at 450 °C type II films. The inset diffraction pattern for type I film shows two sets of hcp rings with hk.0 rings. The set with smaller diameter is for Ti and with larger diameter is for CoCrPt films. The diffraction pattern (inset) for type II films shows two sets of hcp rings with hk.0 indexing representing Ti and CoCrPt and a third set of
bcc rings arising from the CrMn top layer. The presence of only the hk.0 hcp rings in the both types of films shows that they are very strongly 00.1 textured perpendicular films. The grains of the films are small and equiaxed. Both types of the films show similar grain size. No grain growth was observed for annealed films. The mean grain size ($\mu$) of both films is 12.6 nm and the standard deviation ($\sigma$) of the grain size distribution is 2 nm. The similar microstructural features before and after annealing suggest that the increase in coercivity observed after annealing is due to local changes in chemistry.

To demonstrate that interdiffusion of CrMn from the overlayer into the magnetic layer occurs, the compositional distribution was examined by electron energy-loss spectroscopy (EELS). Figure 5 shows typical EELS images taken from a type II film annealed at 450 °C for 5 min. The sample for EELS observation was prepared by careful thinning from the bottom. Once the sample became electron transparent, the sample was again carefully thinned from the top to eliminate the CrMn layer, ensuring that the EELS spectrum was collected from an area that had only the CoCrPt layer present. Figure 5(a) is a zero electron energy-loss image. Figures 5(b)–5(d) are the Cr, Mn, and Co mapping of the same image. From the mapping of Cr, Mn, and Co (bright regions) it is clear that Cr and Mn do diffuse down through the grain boundaries from the top layer. This evidence clearly supports the assertion made earlier in the section that the interdiffusion of Cr and Mn into the film through the grain boundaries during annealing decoupled the grains which improves the magnetic properties.

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