

- 1) Consider the kinematics of $e^+e^- \rightarrow \mu^+\mu^-$ scattering in the center-of-mass. The electrons have momentum \mathbf{k} and the muons emerge at an angle θ relative to the electron and positron direction. Evaluate the expressions for s , t and u , retaining both m_e and m_μ .

Suppose that these beams are colliding head-on. With p indicating the four-momenta,

$$p_{e^-} = \begin{bmatrix} \sqrt{m_e^2 + k^2} \\ k \\ 0 \\ 0 \end{bmatrix} \quad p_{e^+} = \begin{bmatrix} \sqrt{m_e^2 + k^2} \\ -k \\ 0 \\ 0 \end{bmatrix} \quad p_{\mu^-} = \begin{bmatrix} \sqrt{m_\mu^2 + k'^2} \\ k' \cos \theta \\ k' \sin \theta \\ 0 \end{bmatrix} \quad p_{\mu^+} = \begin{bmatrix} \sqrt{m_\mu^2 + k'^2} \\ -k' \cos \theta \\ -k' \sin \theta \\ 0 \end{bmatrix}$$

Define:

$$k'^2 \equiv m_e^2 + k^2 - m_\mu^2$$

$$\begin{aligned} s &\equiv (p_{e^+} + p_{e^-})^2 = (p_{\mu^+} + p_{\mu^-})^2 \\ (p_{e^+} + p_{e^-})^2 &= 2m_e^2 + 2p_{e^+} \cdot p_{e^-} = 2m_e^2 + 2(m_e^2 + k^2 + k^2) \\ &= 4m_e^2 + 4k^2 \end{aligned}$$

$$\begin{aligned} t &\equiv (p_{e^+} - p_{\mu^+})^2 = (p_{e^-} - p_{\mu^-})^2 \\ (p_{e^+} - p_{\mu^+})^2 &= m_e^2 + m_\mu^2 - 2p_{e^+} p_{\mu^+} \\ &= m_e^2 + m_\mu^2 - 2\left(\sqrt{m_e^2 + k^2}\right)\left(\sqrt{m_\mu^2 + k'^2}\right) + 2(kk' \cos \theta) \\ &= m_e^2 + m_\mu^2 - 2(m_e^2 + k^2) + 2(kk' \cos \theta) \\ &= m_\mu^2 - m_e^2 - 2k^2 + 2(kk' \cos \theta) \end{aligned}$$

Simplification was possible since there is no reason to believe the muon energies should differ from the electron energies in this frame.

k' is defined above.

$$\begin{aligned} u &\equiv (p_{e^+} - p_{\mu^-})^2 \\ &= m_e^2 + m_\mu^2 - 2p_{e^+} p_{\mu^-} = m_e^2 + m_\mu^2 - 2\left(\sqrt{m_e^2 + k^2}\right)\left(\sqrt{m_\mu^2 + k'^2}\right) - 2(kk' \cos \theta) \\ &= m_\mu^2 - m_e^2 - 2k^2 - 2(kk' \cos \theta) \end{aligned}$$

k' is defined above.

2) This problem concerns fixed target collisions of protons or anti-proton beams with a proton target.

a) What is the expression for s in terms of the beam proton energy E and the mass m_p ?

$$s = \left(\left[\frac{E}{\sqrt{E^2 - m_p^2}} \right] + \left[\begin{matrix} m_p \\ 0 \end{matrix} \right] \right)^2 = 2m_p^2 + 2Em_p$$

b) What is the heaviest new particle that can be pair-produced (i.e., $pp \rightarrow ppX\bar{X}$) in the collision of an 800 GeV proton with a proton at rest? (Don't forget the protons in the final state!)

At the threshold, \sqrt{s} is the available energy in the CM frame:

$$\begin{aligned} \sqrt{s} &= \sqrt{2m_p^2 + 2 \cdot 800 \text{ GeV} \cdot m_p} = 2m_p + 2m_x = \sqrt{2(938.28)^2 + 2 \cdot (800000) \cdot (938.28)} = 38769 \text{ MeV} \\ \frac{38769 \text{ MeV} - 2m_p}{2} &= m_x = 18446 \text{ MeV} \end{aligned}$$

Looking at a list of particles available from Wikipedia, it looks like one could well get a bottom quark from this interaction. It doesn't look like one could pair-produce any of the bosons predicted by the standard model. It seems unfortunate that despite all of this energy, only a small portion is available in the resulting CM frame for pair production. If they wanted something bigger, they should try a heavier target like some Lead.

c) Repeat b but with $p\bar{p} \rightarrow X\bar{X}$. Does this answer make sense?

$$\begin{aligned} \sqrt{s} &= 2m_x \\ m_x &= \frac{38769 \text{ MeV}}{2} = 19385 \text{ MeV} \end{aligned}$$

This makes perfect sense—there is more available energy in the CM frame.

Again, it looks like the heaviest particle possible is a bottom quark.

3) An electron of energy $E = 4.88 \text{ GeV}$ scatters at $\theta = 10^\circ$ from a proton of mass M at rest. The scattered electron has energy E' and the proton is inelastically excited to a state of invariant mass W .

a) Calculate E' in terms of E, W, M, θ .

Let m_e be negligibly small compared to E .

Then:

$$\left(\begin{bmatrix} E \\ E \\ 0 \end{bmatrix} - \begin{bmatrix} E' \\ E' \cos \theta \\ E' \sin \theta \end{bmatrix} \right)^2 = \left(\begin{bmatrix} M \\ 0 \\ 0 \end{bmatrix} - \begin{bmatrix} E_N \\ E - E' \cos \theta \\ -E' \sin \theta \end{bmatrix} \right)^2$$

$$2m_e^2 - 2EE'(1 - \cos \theta) = M^2 + W^2 - 2ME_N$$

$$\frac{2m_e^2}{2M} - E' \frac{E}{M} (1 - \cos \theta) = \frac{M^2 + W^2}{2M} - E_N$$

$$\frac{m_e}{M} \approx 0$$

$$\therefore E' \frac{E}{M} (1 - \cos \theta) = E_N - \frac{M^2 + W^2}{2M}$$

Energy conservation: $M + E = E_N + E'$

$$E' \frac{E}{M} (1 - \cos \theta) = M + E - E' - \frac{M^2 + W^2}{2M}$$

$$E' \left(1 + \frac{E}{M} (1 - \cos \theta) \right) = M + E - \frac{M^2 + W^2}{2M}$$

$$E' = \frac{E + \frac{M^2 - W^2}{2M}}{1 + \frac{E}{M} (1 - \cos \theta)}$$

b) What is the answer when $W = M$? Does this make sense?

$$E' = \frac{E + \frac{M^2 - W^2}{2M}}{1 + \frac{E}{M} (1 - \cos \theta)} \rightarrow \frac{E}{1 + \frac{E}{M} (1 - \cos \theta)}$$

This is just the Compton scattering formula—for elastic scattering (whereas the nucleus has not picked up any bonus mass).

c) Use this formula to calculate E' for the cases of $W = M$, $W = 1230\text{MeV}$, and $W = 1520\text{MeV}$.

$$E'(W = M) = \frac{4880 \text{ MeV}}{1 + \frac{4880 \text{ MeV}}{938.3 \text{ MeV}}(1 - \cos 10^\circ)} = 4523 \text{ MeV}$$

$$E'(W = 1230 \text{ MeV}) = \frac{4880 \text{ MeV} + \frac{(938.3 \text{ MeV})^2 - (1230 \text{ MeV})^2}{2 \cdot 938.3 \text{ MeV}}}{1 + \frac{4880 \text{ MeV}}{938.3 \text{ MeV}}(1 - \cos 10^\circ)} = 4211 \text{ MeV}$$

$$E'(W = 1520 \text{ MeV}) = \frac{4880 \text{ MeV} + \frac{(938.3 \text{ MeV})^2 - (1520 \text{ MeV})^2}{2 \cdot 938.3 \text{ MeV}}}{1 + \frac{4880 \text{ MeV}}{938.3 \text{ MeV}}(1 - \cos 10^\circ)} = 3816 \text{ MeV}$$

d) Consider the limit from part (a) where $W - M \ll M$. Explain how this limit allows us to easily obtain energy levels of the Oxygen nucleus from the data in Povh 6.3.

$$E' = \frac{E + \frac{M^2 - W^2}{2M}}{1 + \frac{E}{M}(1 - \cos \theta)} \rightarrow \frac{E + \frac{(M+W)(M-W)}{2M}}{1 + \frac{E}{M}(1 - \cos \theta)}$$

$$\alpha = W - M$$

$$\frac{(M+W)(M-W)}{2M} \rightarrow -\frac{2M\alpha + \alpha^2}{2M} \approx -\alpha$$

$$\therefore E'_{W-M \ll M} \approx \frac{E - (W - M)}{1 + \frac{E}{M}(1 - \cos \theta)} =$$

$$W_{W-M \ll M} = E + M - E' \left[1 + \frac{E}{M}(1 - \cos \theta) \right]$$

From Povh 6.3, it is given that $E = 246 \text{ MeV}$, $\theta = 148.5^\circ$. The rest mass of an Oxygen-16 nucleus is $\approx 8M_H + 8M_n = 8 \cdot 938.78 \text{ MeV} + 8 \cdot 939.57 \text{ MeV} = 15027 \text{ MeV}$.

Of course, this approximation seems unnecessary: Instead, I can take the equally benign and more accurate formula from the original expression

$$W^2 = M^2 - 2M \left(\left[1 + \frac{E}{M}(1 - \cos \theta) \right] E' - E \right)$$

All that's left, then, is to read off the E' at the spikes to find the excited states. The big one, for example, corresponds to roughly $E = 160$ for

$$W_1 \approx 15108 \text{ MeV} .$$

- 4) When thinking about the number of independent parameters in deep inelastic scattering, we decided that there were two independent variables which we chose as q^2 and W . That means that any other scalar quantities can be expressed in terms of these. Taking the electron 4-momentum before (after) scattering to be $p(p')$, with $q = p - p'$, and the 'target system' before (after) to be $P(P')$, express the following in terms of q^2 , W , and the proton mass M . We continue to neglect m_e .

- $P \cdot q$
- $P' \cdot q$
- $p \cdot p'$

(fill in derivation of first one).

$$W^2 = (P')^2 = (P + q)^2 = M^2 + q^2 + 2P \cdot q$$

$$\therefore P \cdot q = \frac{W^2 - M^2 - q^2}{2}$$

Next,

$$p + P = p' + P'$$

$$P + p - p' = P'$$

$$P + q = P'$$

$$(P + q) \cdot q = P' \cdot q$$

$$\therefore P' \cdot q = P \cdot q + q^2 = \frac{W^2 - M^2 + q^2}{2}$$

Finally,

$$-\frac{1}{2}q^2 = -(p - p') \cdot (p - p') = -\frac{1}{2}(p^2 + p'^2 - 2p \cdot p') = -\frac{1}{2}m_e^2 + p \cdot p' \approx p \cdot p'$$